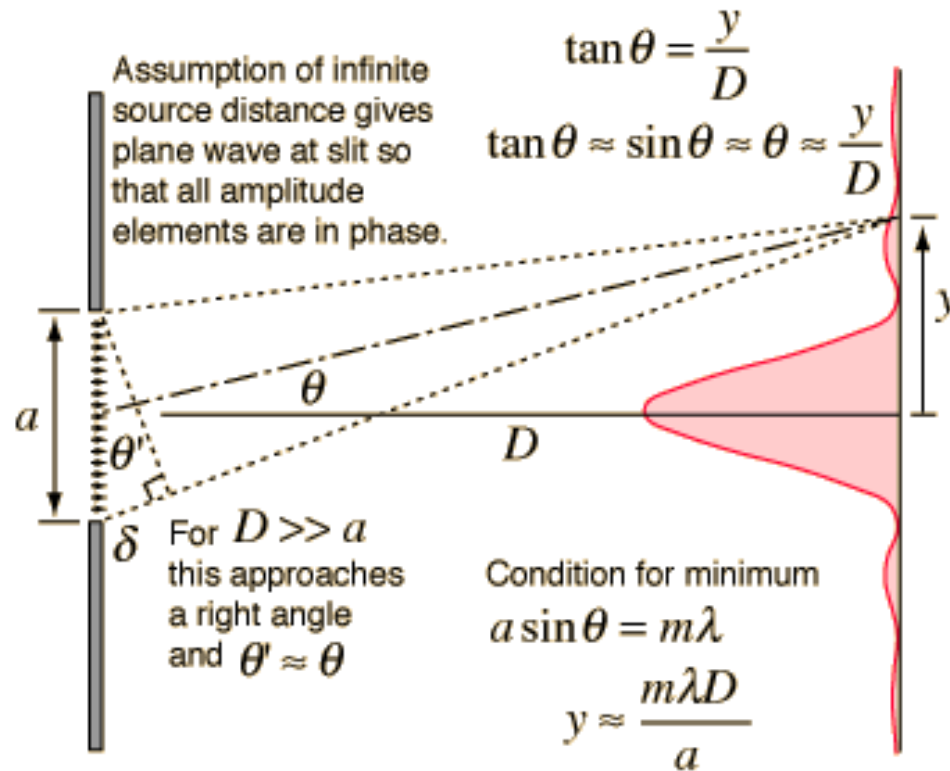


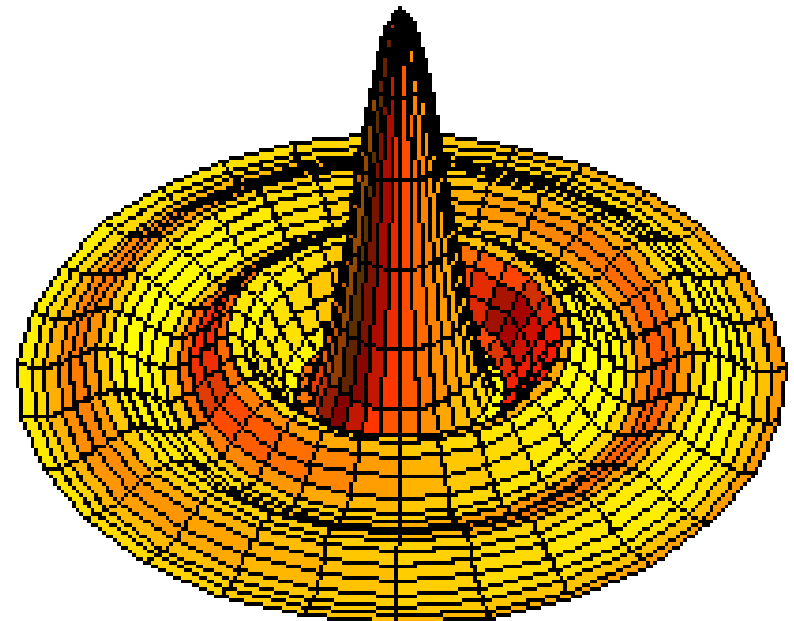
Optical Resolution



Fraunhofer Diffraction

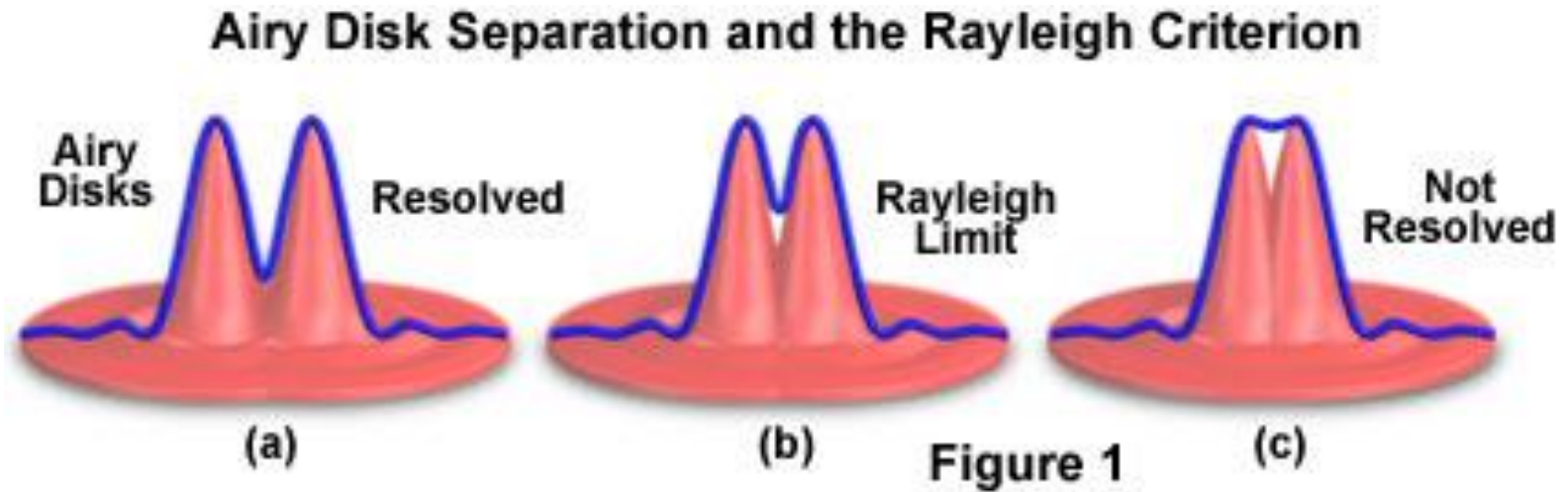
Single slit & circular aperture

Circular aperture -> Airy disk



Airy Disk

Optical Resolution



$$d = \frac{0.61\lambda}{NA}$$

d = Rayleigh Criterion length

λ = wavelength

NA : Objective Numerical Aperture

For visible wavelength

$d \sim 250\text{nm}$

Fast, background-free, 3D super-resolution optical fluctuation imaging (SOFI)

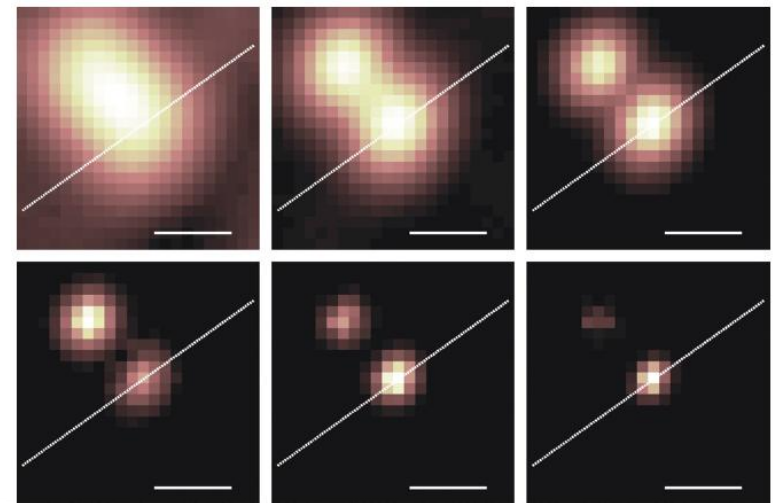
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Departments of ^aChemistry and Biochemistry and ^bPhysiology, and ^cCalifornia NanoSystems Institute, University of California, Los Angeles, CA 90095; and ^dIII. Institute for Physics, Georg-August-University, 37073 Göttingen, Germany

Edited by John W. Sedat, University of California, San Francisco, CA, and approved October 29, 2009 (received for review July 15, 2009)

Super-resolution optical microscopy is a rapidly evolving area of fluorescence microscopy with a tremendous potential for impacting many fields of science. Several super-resolution methods have been developed over the last decade, all capable of overcoming the fundamental diffraction limit of light. We present here an approach for obtaining subdiffraction limit optical resolution in all three dimensions. This method relies on higher-order statistical analysis of temporal fluctuations (caused by fluorescence blinking/intermittency) recorded in a sequence of images (movie). We demonstrate a 5-fold improvement in spatial resolution by using a conventional wide-field microscope. This resolution enhancement is achieved in iterative discrete steps, which in turn allows the evaluation of images at different resolution levels. Even at the lowest level of resolution enhancement, our method features significant background reduction and thus contrast enhancement and is demonstrated on quantum dot-labeled microtubules of fibroblast cells.

cumulants | fluorescence | quantum dots | superresolution microscopy | intermittency



Statistical order) while simultaneously enhancing the image contrast. We argue that no other super-resolution microscopy technique can compete with the simplicity of the SOFI approach and its undemanding requirements with regard to fluorescent labels, optics, and other hardware. The experimental procedure essentially amounts to taking a movie of a fluctuating signal.

Seoncheol Cha

Department of Physics, Sogang University

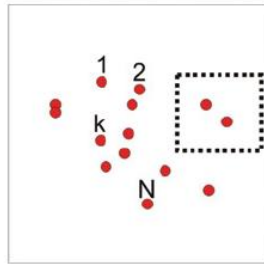
Soft Matter Optical Spectroscopy Laboratory

Principle of SOFI

Emitter distribution

In the object plane

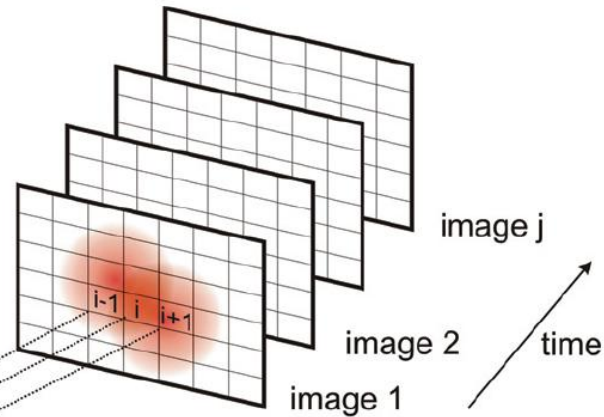
A



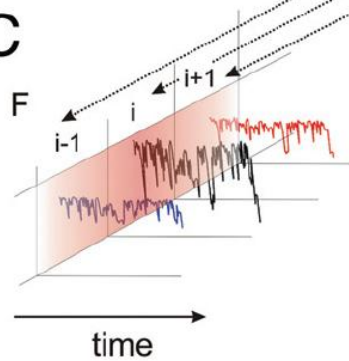
Magnified detail of A

: Conventional imaging cannot resolve two points

B

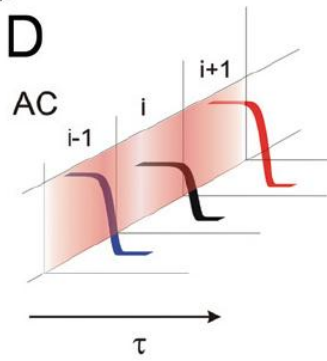


C



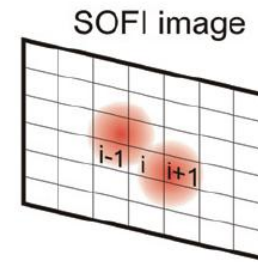
Time trace of B

D



2nd order
correlation function of C

E



Two points are
Resolved!

Three Condition for SOFI

1. The fluorescent label has to exhibit **at least two different emission states**. For example, these states can be a fluorescent and a nonfluorescent one, but in principle any two or more states that are optically distinguishable will do.
2. Different emitters have to switch between states **repeatedly** and independently from each other in a **stochastic** way.
3. For this approach, the image should be acquired with **pixels smaller than the diffraction limit**. Resolution less than the pixel size will be the topic of a future publication.

Theory 1 – Correlation Function

Fluorescence source distribution

$$\sum_{k=1}^N \delta(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_k \cdot s_k(t)$$

ε_k : constant molecular brightness
 $s_k(t)$: time-dependent fluctuation

Fluorescence signal at position \mathbf{r} and time t

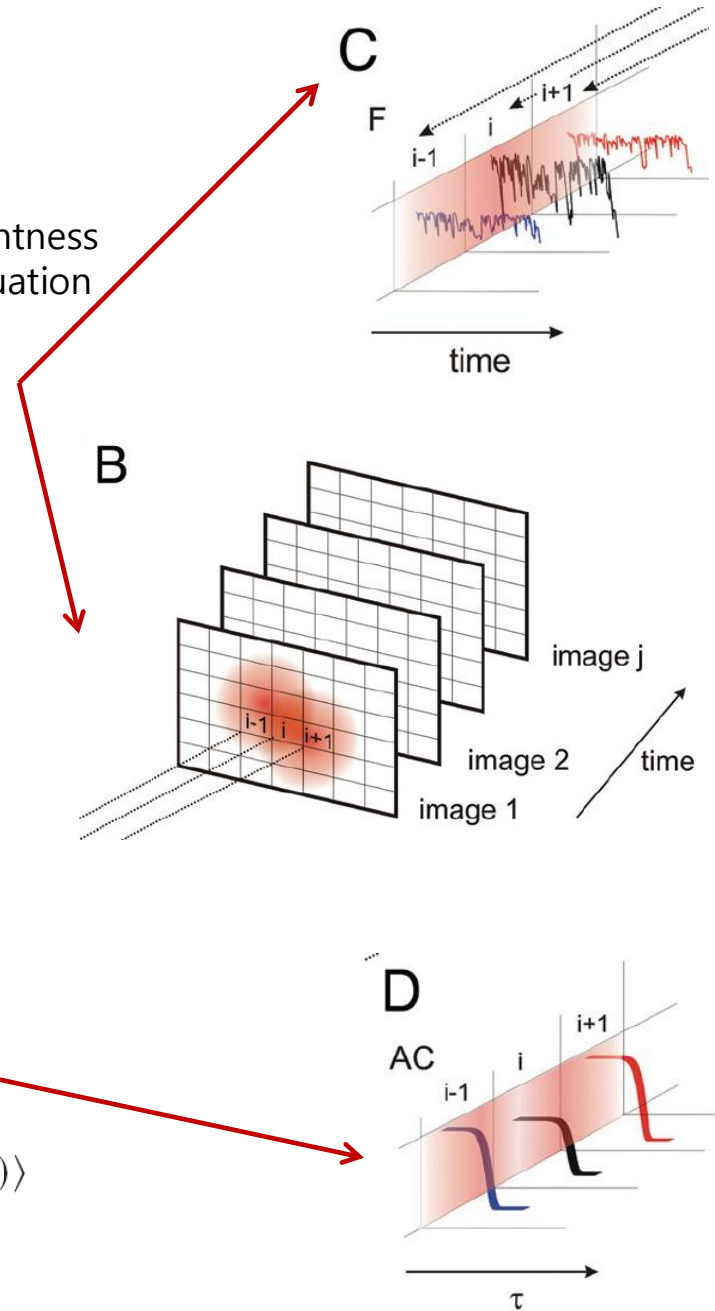
$$F(\mathbf{r}, t) = \sum_{k=1}^N U(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_k \cdot s_k(t)$$

Fluorescence Fluctuation

$$\begin{aligned} \delta F(\mathbf{r}, t) &= F(\mathbf{r}, t) - \langle F(\mathbf{r}, t) \rangle_t \\ &= \sum_k U(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_k \cdot [s_k(t) - \langle s_k(t) \rangle_t] \\ &= \sum_k U(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_k \cdot \delta s_k(t), \end{aligned}$$

2nd order correlation function

$$\begin{aligned} G_2(\mathbf{r}, \tau) &= \langle \delta F(\mathbf{r}, t + \tau) \cdot \delta F(\mathbf{r}, t) \rangle_t \\ &= \sum_{j,k} U(\mathbf{r} - \mathbf{r}_j) U(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_j \cdot \varepsilon_k \cdot \langle \delta s_j(t + \tau) \delta s_k(t) \rangle \\ &= \sum_k U^2(\mathbf{r} - \mathbf{r}_k) \cdot \varepsilon_k^2 \cdot \langle \delta s_k(t + \tau) s_k(t) \rangle \end{aligned}$$



Theory 2 – Correlation Function to Cumulant Function

PSF in 2nd order Image

$$U(\mathbf{r}) = \exp\left(-\frac{x^2 + y^2}{2\omega_0^2} - \frac{z^2}{2\omega_{z0}^2}\right) \Rightarrow U^2(\mathbf{r}) = \exp\left(-\frac{x^2 + y^2}{2\tilde{\omega}_0^2} - \frac{z^2}{2\tilde{\omega}_{0z}^2}\right)$$

Generalize correlation function

$$\tilde{\omega}_{0z} = \omega_{0z}/\sqrt{2} \text{ and } \tilde{\omega}_0 = \omega_0/\sqrt{2}$$

$$G_n(\mathbf{r}, \tau_1, \dots, \tau_{n-1})$$

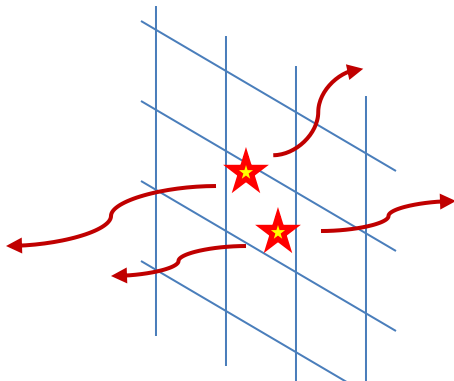
$$= \langle \delta F(\mathbf{r}, t) \delta F(\mathbf{r}, t + \tau_1) \cdots \delta F(\mathbf{r}, t + \tau_{n-1}) \rangle_t$$

Cumulant function $C_n(\mathbf{r}, \tau_1, \dots, \tau_{n-1})$

$$C_n(\mathbf{r}, \tau_1, \dots, \tau_{n-1}) = \sum_k U^n(\mathbf{r} - \mathbf{r}_k) \varepsilon_k^n w_k(\tau_1, \dots, \tau_{n-1})$$

w_k : correlation-based weighting function

All cross terms caused by lower-order correlation contributions are eliminated in cumulants.



Consider 4 photons from two emitters in one pixel,

4th order correlation contribution is concealed by lower order (2nd order) contribution

Theory 3 – Cumulant Function

$$C_n(\mathbf{r}, \tau_1, \dots, \tau_{n-1}) = \sum_k U^n(\mathbf{r} - \mathbf{r}_k) \varepsilon_k^n w_k(\tau_1, \dots, \tau_{n-1})$$

Resolution is enhanced by a factor \sqrt{n}
n-fold larger brightness will appear 2^n times brighter in n^{th} order Image

$$C_2(\mathbf{r}, \tau_1) = G_2(\mathbf{r}, \tau_1)$$

$$C_3(\mathbf{r}, \tau_1, \tau_2) = G_3(\mathbf{r}, \tau_1, \tau_2)$$

$$\begin{aligned} C_4(\mathbf{r}, \tau_1, \tau_2, \tau_3) = & G_4(\mathbf{r}, \tau_1, \tau_2, \tau_3) - G_2(\mathbf{r}, \tau_1) \cdot G_2(\mathbf{r}, \tau_3) \\ & - G_2(\mathbf{r}, \tau_1 + \tau_2) \cdot G_2(\mathbf{r}, \tau_2 + \tau_3) - G_2(\mathbf{r}, \tau_1 + \tau_2 + \tau_3) \cdot G_2(\mathbf{r}, \tau_2) \end{aligned}$$

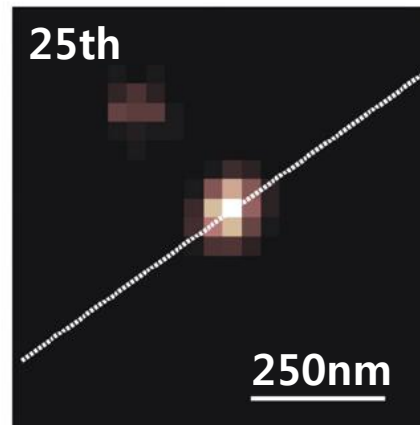
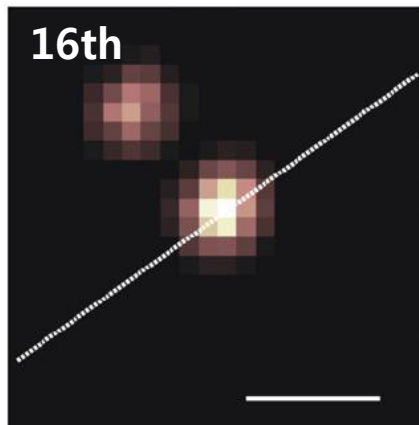
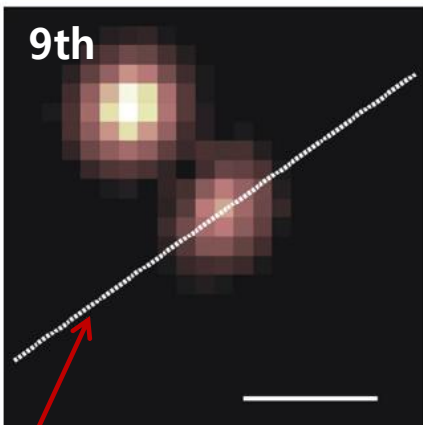
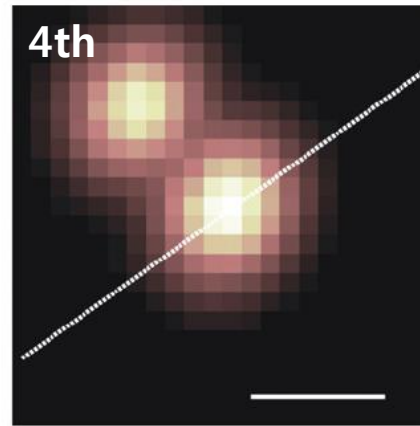
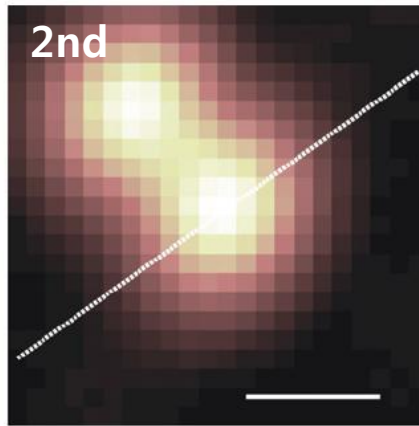
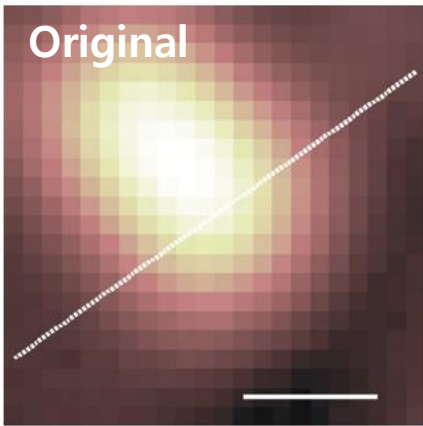
Setup

Microscope Setup and Data Analysis. Movies were taken on an inverted wide-field microscope (Olympus IX71, objective: UPlanApo 60 \times , 1.45, Oil). A 470-nm LED array device was used as a light source (Lumencor Aura Light Engine) and a CCD camera (Andor iXon + 885) was used to record the signal. Filter set was dichroic [505 DCXR, Chroma Technology; emission (D620/40, Chroma Technology)]. Magnification was adjusted to obtain 35 nm per pixel. To generate and evaluate SOFI images, movies were analyzed by using a custom-written Matlab (Mathworks) code. The shortest accessible time lag is the frame integration time (time between two subsequent frames). We computed all SOFI images for the zero time lag only $C_n(\mathbf{r}, 0, 0, \dots, 0)$. In this case a computationally less expensive expression for the cumulants formula can be used (see [SI Text](#)). Cumulants of orders > 2 can turn negative depending on the underlying fluctuation pattern. SOFI images are therefore displayed as absolute values.

Results 1 – Images of Two Quantum Dots due to Correlation Orders

S/B : 1.9
FWHM : 289nm

S/B : 250

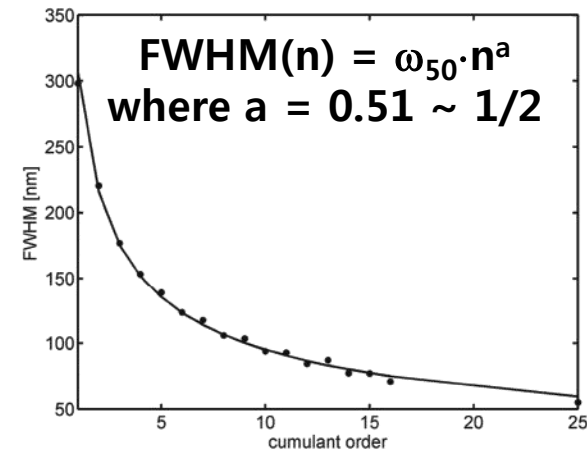
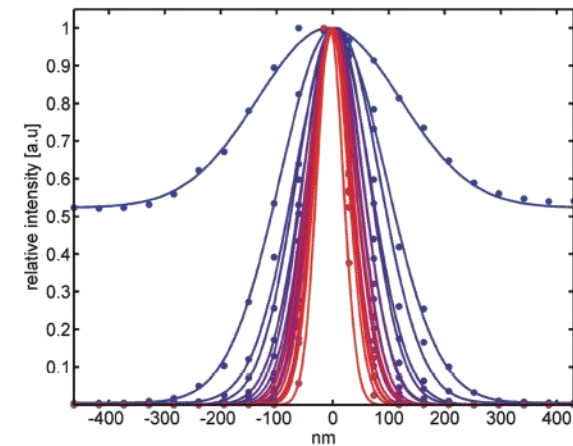


FWHM : 55nm

Cross-section of
Original Image

Signal to Background Ratio (S/B)
2,000 frames (100 ms per frame)

Profiles & Fits



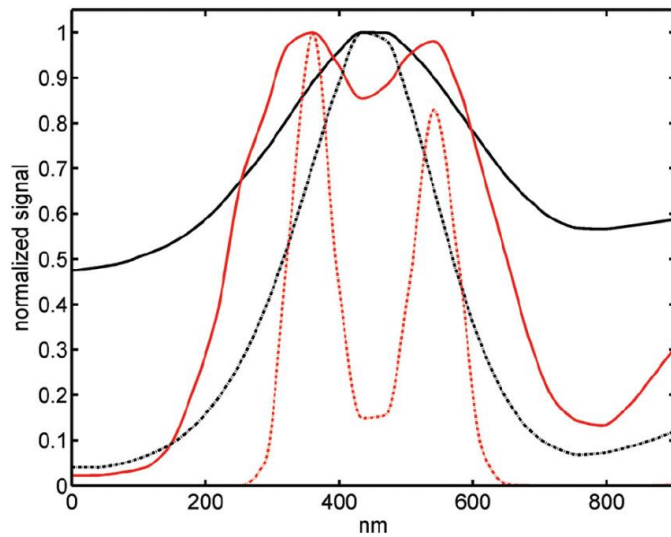
**FWHM versus
Correlation Order**

Results 2 – Point Spread Function from Single Quantum Dot

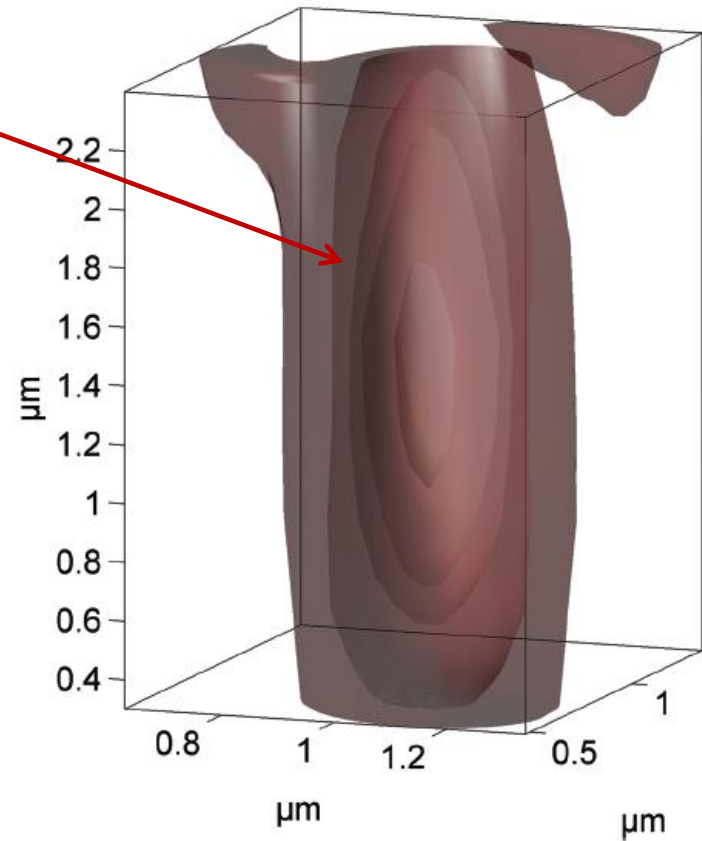
Original, 2nd, 4th, 9th and 16th order Image

1/e² iso-surfaces

Results 3 – Resolution Enhancement



- Original Image
- 2nd order Image
- Cross-section of Original Image
- Cross-section of 2nd order Image

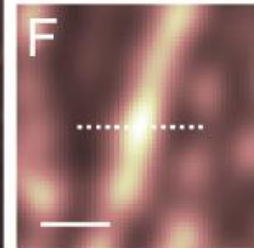
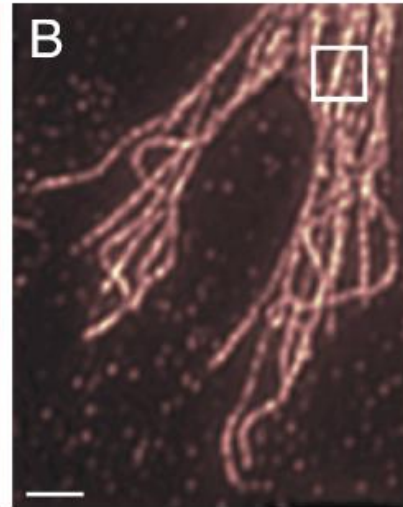
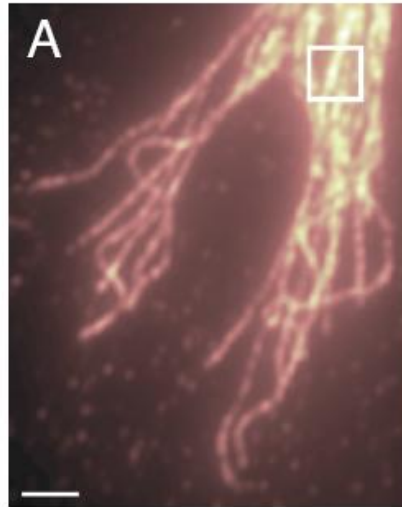


300nm spacing scan image
4,000 frames (75 ms per frame)

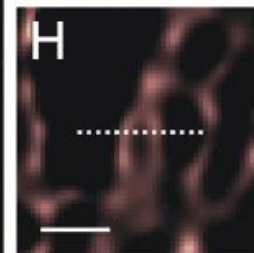
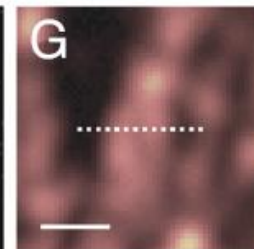
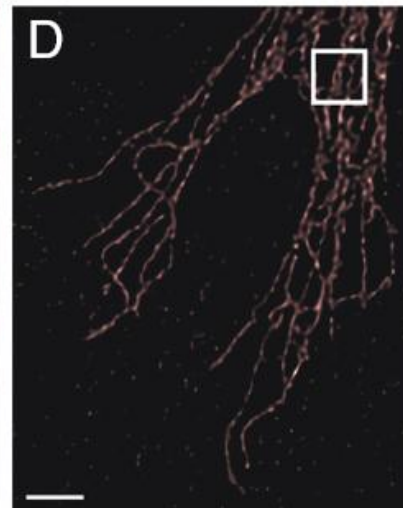
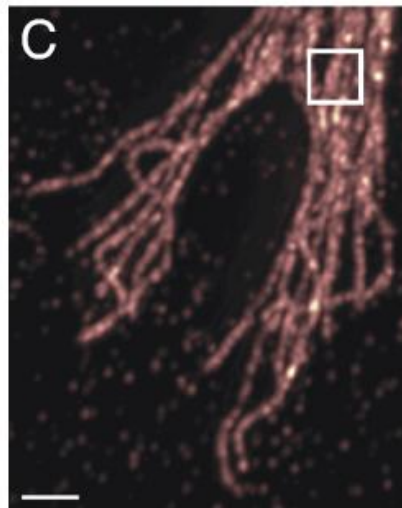
Results 4 – SOFI Images of Cells (α -tubulin network of a 3T3 fibroblast cell)

A deconvolved

Original Image



2nd order Image

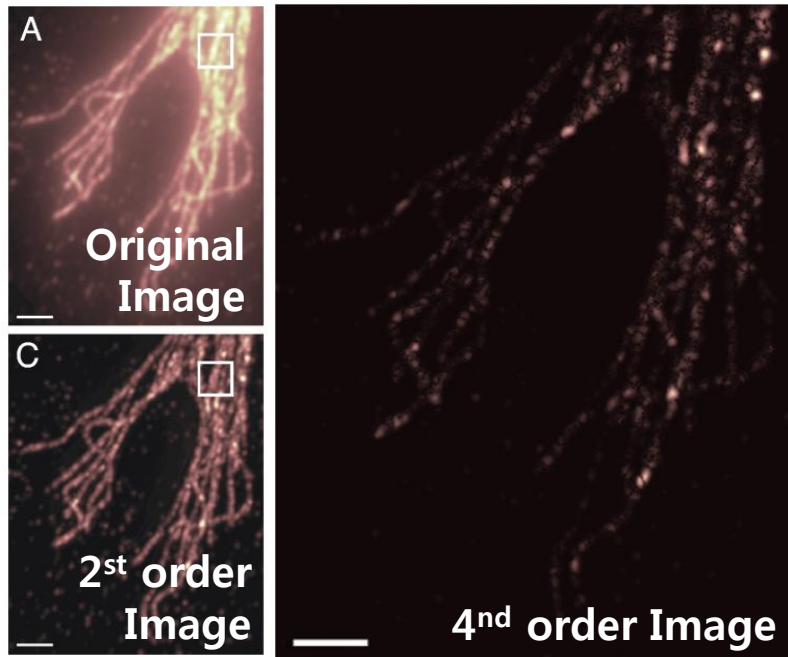


C deconvolved

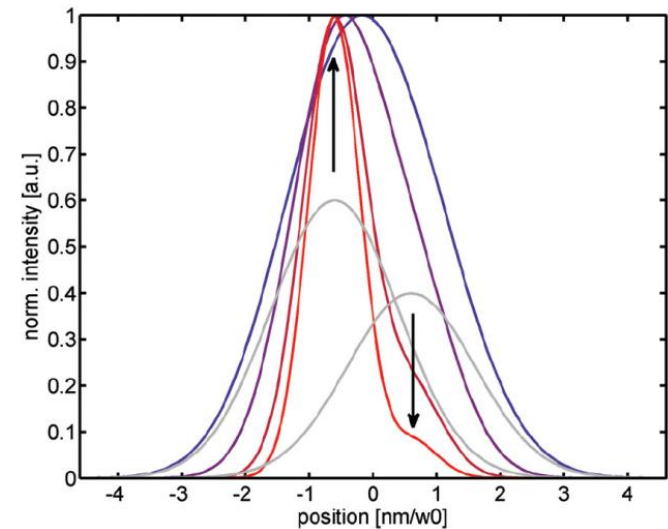
QD625 Labeled
3,000 frames (100 ms per frame)

Limitation 1. The Brightness Scaling of the Images

- n-fold larger brightness will appear 2^n times brighter in n^{th} order Image
- > very high Dynamic Range
 - > Masking effect of dim emitters in proximity to bright emitters



Blue to red : cumulant order 1,2,4,6
Gray : individual diffraction limit



Blinking rate : 10% of Frame rate

Limitation 2. Limited measurement time

Theoretically, SOFI eliminate any kinds of noise

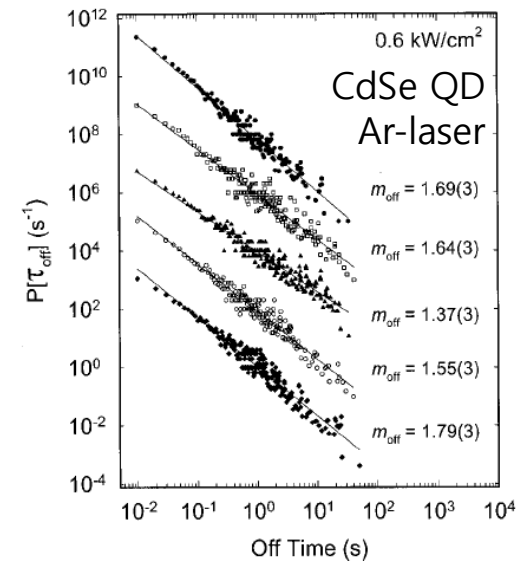
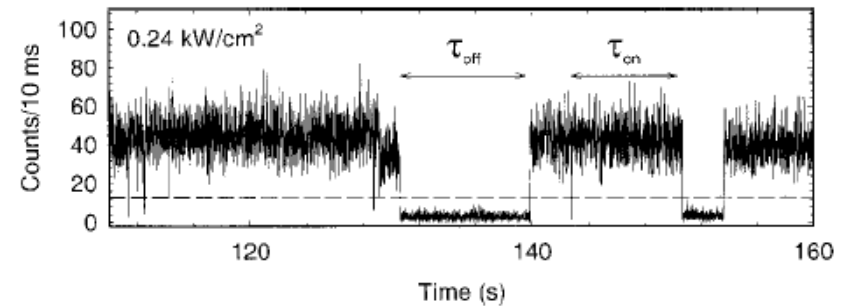
Acquisition times are limited
(by photobleaching ...)

Potential to super-resolution imaging at high frame rates

QDs remain in the on state for a few seconds leading to very different brightness values in image

-> Need to uniform and fast blinking emitter for SOFI at high frame rates

Kuno *et al.*, J. Chem. Phys (2001)
Power law on/off distribution



Conclusion

statistical order) while simultaneously enhancing the image contrast. We argue that no other super-resolution microscopy technique can compete with the simplicity of the SOFI approach and its undemanding requirements with regard to fluorescent labels, optics, and other hardware. The experimental procedure essentially amounts to taking a movie of a fluctuating signal.

ground reduction. Last, SOFI is not limited to blinking between fluorescent on and off states. Any (even nonfluorescent) fluctuating objects, such as rotating dipoles, or blinking of celestial objects, such as binary stars, could be imaged and superresolved by SOFI.