Infrared Spectroscopy of Ionic Liquids: Quantitative Aspects and Determination of Optical Constants

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The infrared (IR) spectra of ionic liquids involving 1-butyl-3-methylimidazolium (BMI) and the (CF3SO2)2N, BF4, or PF6 anions, recorded in the **transmission and attenuated total reflection (ATR) modes, exhibit strong differences in the most intense anion absorption profiles**. These distortions come from optical effects and make difficult any quantitative analysis of, for example, the antisymmetric stretching vibrations of the BF4 and PF6 anions. A method is proposed to determine the optical constants, from which any type of experimental spectrum can be simulated. It is then possible to use the anion vibrational bands as spectroscopic probes of the local interactions occurring in the neat ionic liquids and in solutions. This is illustrated by a direct identification of ion pairs and separated ions in the IR spectra of BMI–PF6 solutions.

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IR Spectroscopy for Ionic Liquids

conformational state of cations & anions (with Raman spectroscopy & *ab initio* calculations)

interactions within ion pairs & larger clusters

Introduction

IR spectra of ionic liquids exhibit intense anion absorption bands (high concentration)

 $<0.5 \ \mu m$ thickness are required in the FT-IR

Comparison of the Infrared Spectra of Ionic Liquids Obtained Under Different Conditions [BMIM][PF6]



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Wavenumber/cm⁻¹

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Fresnel Reflection coefficients

 $\hat{r}_{v} = R_{v}^{\frac{1}{2}} \exp(i\theta_{v}),$ $\hat{r}_{s} = \frac{n_{0} \cos \alpha_{0} - \hat{n}_{t} \cos \alpha_{s}}{n_{0} \cos \alpha_{0} + \hat{n}_{t} \cos \alpha_{s}}$ $\hat{r}_{p} = \frac{\hat{n}_{t} \cos \alpha_{0} - n_{0} \cos \alpha_{p}}{\hat{n}_{t} \cos \alpha_{0} + n_{0} \cos \alpha_{p}}$

Snell's law

 $n_0 \sin \alpha_0 = \hat{n}_t \sin \alpha_s = \hat{n}_n \sin \alpha_p$

 $\hat{\varepsilon} = \hat{n}^2$

$$\hat{\epsilon}_t / \epsilon_0 = \sin^2 \alpha_0 + [(1 - \hat{r}_s) / (1 + \hat{r}_s)]^2 \cos^2 \alpha_0$$

$$\hat{\varepsilon}_n/\varepsilon_0 = \sin^2 \alpha_0 / [1 - (\hat{\varepsilon}_t/\varepsilon_0) \{ (1 - \hat{r}_p) / (1 + \hat{r}_p) \}^2 \cos^2 \alpha_0]$$

 $\theta_s' + \Delta \theta_s \simeq \theta_s$

M.J.Dignam&S.Mamiche-Afara, Spectrochimica Acta (1988)

Fresnel Reflection coefficients

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Snell's law

 $n_0 \sin \alpha_0 = \hat{n}_t \sin \alpha_s = \hat{n}_n \sin \alpha_p$

Kramers – Kronig relation with F.T.

$$\theta'_{v}(\bar{v}) = -4 \int_{0}^{\infty} d\delta \sin(2\pi \bar{v}\delta)$$
$$\times \int_{0}^{\infty} d\bar{v}_{0} \ln [R_{v}(\bar{v}_{0})]^{\frac{1}{2}} \cos(2\pi \bar{v}_{0}\delta)$$

Phase correction

$$\Delta \theta_s = \theta_s - \theta'_s = 2 \arctan\left[(n_0^2 \sin^2 \alpha_0 - n_t^2)^{\frac{1}{2}}/n_0 \cos \alpha_0\right] - \theta'_s,$$

$$\Delta \theta_p = \theta_p - \theta'_p = 2 \arctan\left[(n_0^2 \sin^2 \alpha_0 - n_n^2)^{\frac{1}{2}}n_0/n_n n_t \cos \alpha_0\right] - \theta'_p,$$

M.J.Dignam&S.Mamiche-Afara, Spectrochimica Acta (1988)

$$\delta(\mathbf{v}_{\mathrm{a}}) = I + \frac{2}{\pi} P \int_0^\infty \frac{\nu . \ln \sqrt{R_{\mathrm{s}}(\nu)}}{\nu^2 - \nu_{\mathrm{a}}^2} \mathrm{d}\nu$$

$$I = 2 \arctan \frac{\sqrt{n_0^2 \sin^2 \theta_0 - n^2(v_u)}}{n_0 \cos \theta_0}$$

$$n(\mathbf{v}) = n_0 \text{Re} \left[\sin^2 \theta_0 + \left[\frac{1 - \hat{r}_s(\mathbf{v})}{1 + \hat{r}_s(\mathbf{v})} \right]^2 \cos^2 \theta_0 \right]^{1/2}$$
$$k(\mathbf{v}) = -n_0 \text{Im} \left[\sin^2 \theta_0 + \left[\frac{1 - \hat{r}_s(\mathbf{v})}{1 + \hat{r}_s(\mathbf{v})} \right]^2 \cos^2 \theta_0 \right]^{1/2}$$

For K-K relation, n(v, large) ~ constant k(v, large) ~ small

n(v, small) ~ neglect k(v, small) ~ neglect Determination of the Optical Constants of Ionic Liquids from the ATR Spectra



Simulation of Various Infrared Spectra from the Optical Constants



Pure Ionic Liquids



TABLE I. Observed anion vibrations for BMI-PF₆ at room temperature.

TABLE II.	Observed anion	vibrations fo	r BMI-BF ₄ at	room temperature.

BMI-PF ₆	v_1	V ₂	V ₃
O _h	A _{1g} , R	E _g , R	F _{1u} , IR
C_{4v}	A1, IR, R	A_1 , IR, R, B_1 , R	A1, IR, R, E, IR, R
IR, exp cm ⁻¹	741 ^a		841
Raman, exp cm ⁻¹	741	567	864 ^a

^a Vibrations normally IR or Raman inactive in the O_h symmetry.

BMI-BF ₄	ν_1	V2	V3
	A_1, R A_2, R, R	E, R E IR R	T ₂ , IR, R F. IR, R, A, IR, R
\mathbb{R} , exp cm ⁻¹	763 ^a	2, 11, 11	1061, 1039
Raman, exp cm ⁻¹	764.6	353	1056

^a Vibrations normally IR inactive in the T_d symmetry.

From Pure Ionic Liquids to Dilute Solutions



BMI-PF ₆	ν_1	V ₂	
	A _{1g} , R	E _g , R	
v	A_1 , IR, R	A_1 , IR, R, B_1 , R	A1, 1

TABLE I. Observed anion vibrations for BMI-PF₆ at room temperature.

V3

F_{1u}, IR

A1, IR, R, E, IR, R

841 864^a

^a Vibrations normally IR or Raman inactive in the O_h symmetry.

741^a

741

High dielectric constant Small v1 IR signal Oh > C4v

 O_h

 C_{4v} IR, exp cm⁻¹

Raman, $exp cm^{-1}$

- A : pure
- B : DMC, dymethyl carbonate (ϵ =3.1)

567

- C : Ac, acetone (ϵ =20.7)
- D : PC, propylene carbonate (ϵ =64.4)
- E : CAN-d₃, deuterated acetonitrile (ϵ =35)

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Conclusion

Careful analysis is required in ATR measurement (especially, low refractive index crystal and strong absorption)

Optical constants could be obtained by IR measurement.

- -> ultimate quantitative information
- -> oscillator strength effects, anisotropic orientations, simulating IR spectra

Anion vibrations constitute a good spectroscopic probe

Imidazolium-Based Ionic Liquids: Quantitative Aspects in the Far-Infrared Region

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The optical constants of some imidazolium-based ionic liquids (ILs) are determined in the mid- and farinfrared regions by combining polarized attenuated total reflection (ATR) and transmittance spectra. The internal vibrations of the cations and anions and the interionic vibrations can thus be quantitatively evaluated. A comparison of the far-IR spectral response of several imidazolium derivatives associated with the $(CF_3SO_2)_2N^-$ anion shows that methylation of the more acidic $C_{(2)}H$ imidazolium group does not change the far-IR intensity and hence that the CH···anion hydrogen bonds play a negligible role compared with electrostatic interactions. The calculated spectra of ion-pair dimers reproduce the far-IR density of states better than those of simple ion pairs.

T.Buffeteau, J.Grondin, Y.Danten, J.C.Lassegues, JPCB (2010)

$$\delta(\mathbf{v}_{\mathrm{a}}) = I + \frac{2}{\pi} P \int_{0}^{\infty} \frac{\mathbf{v} \cdot \ln \sqrt{R_{\mathrm{s}}(\mathbf{v})}}{\mathbf{v}^{2} - \mathbf{v}_{\mathrm{a}}^{2}} \mathrm{d}\mathbf{v}$$

$$I = 2 \arctan \frac{\sqrt{n_0^2 \sin^2 \theta_0 - n^2(v_u)}}{n_0 \cos \theta_0}$$

$$n(v) = n_0 \text{Re} \left[\sin^2 \theta_0 + \left[\frac{1 - \hat{r}_s(v)}{1 + \hat{r}_s(v)} \right]^2 \cos^2 \theta_0 \right]^{1/2}$$

$$k(\mathbf{v}) = -n_0 \text{Im} \left[\sin^2 \theta_0 + \left[\frac{1 - \hat{r}_s(\mathbf{v})}{1 + \hat{r}_s(\mathbf{v})} \right]^2 \cos^2 \theta_0 \right]^{1/2}$$



For K-K relation, n(v, large) ~ constant k(v, large) ~ small

n(v, small) ~ neglect k(v, small) ~ neglect a



Comparison of the optical constants



Well-agreed with result from THz measurement