## Measurement of the coherence time of the Light from a Quasi-thermal Source

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(Received: February 24, 1996)

The photocount distribution from a quasi-thermal light source, a moving ground glass disk (surface roughness; $9 \mu \mathrm{~m}$ ) illuminated by a well-stabilized $\mathrm{He}-\mathrm{Ne}$ laser,. is measured by a photon counting system, and analyzed with theoretical calculations. The distribution approaches the Poisson distribution for the long coherence time $\tau_{\mathrm{F}}$ compared to the measuring time $T$.

The coherence time $\tau_{1}$ of the quasi-thermal source can be changed by controlling the velocity $v$ of the motor driving the glass disk. By the comparison of experimental results and theory for the condition of $T / \tau_{i} \gg 1$, the coherence time $\tau_{i}$ of the quasi-thermal source is turned out to be in the range of $31.43 \mu \mathrm{~s} \sim 2.48 \mu \mathrm{~s}$ according to the circumferential velocity of the disk, and compared with the simple calculation of $\sigma / v$.

## Chaotic light <br> (thermal cavity, filament lamp)

The different atoms are excited by an electrical discharge and emit their radiation

The shape of an emission line is determined by the statistical spread in atomic velocities and the random occurrence of collisions.
$\Delta \nu \gg \Delta v=\frac{1}{\tau_{c}}$
Generally,


Lorentzian lineshape

## collision-broadened light source

## $\gamma_{511}^{510 \mathrm{keV}}$



Gaussian lineshape

Doppler broadening light source


## Quasi-thermal Source

It has property of chaos light, but coherence time is longer than chaos light.

Used rotating ground glass disk (surface roughness: $9 \mu \mathrm{~m}$ )
$\tau_{c}$ of the quasi-thermal source can be changed by controlling the velocity $v$ of the motor driving the glass disk.

$$
\tau_{c}=\frac{\sigma}{\nu} \quad(\text { surface roughness })
$$

$\bar{I}(t, T)$ is the mean intensity that falls on the phototube during the period from t to $\mathrm{t}+\mathrm{T}$, then,

$$
\bar{I}(t, T)=\frac{1}{T} \int_{t}^{t+T} \bar{I}(t) d t^{\prime}
$$

The probability of detecting n number of photons during time duration T is

$$
\begin{aligned}
& P_{n}(T)=\left\langle P_{n}(t, T)\right\rangle \\
& =\left\langle\frac{[\zeta \bar{I}(t, T) T]^{n}}{n!} \exp [\zeta \bar{I}(t, T) T]\right\rangle
\end{aligned}
$$

This result is Known as the Mandel formula. $\zeta$ (efficiency of detector)

So Mean number of photocounts is

$$
\langle n\rangle=\sum_{n=0}^{\infty} n P_{n}(T)=\langle\zeta \bar{I}(t, T) T\rangle=\zeta \overline{I T}
$$

And the second moment of the distribution given by

$$
\begin{aligned}
& \left\langle n^{2}\right\rangle=\sum_{n=0}^{\infty} n^{2} P_{n}(T)=\langle\zeta \bar{I}(t, T) T\rangle+\left\langle[\zeta \bar{I}(t, T) T]^{2}\right\rangle \\
& =\langle n\rangle+\left\langle[\zeta \bar{I}(t, T) T]^{2}\right\rangle
\end{aligned}
$$

The variance of the photocount distribution is therefore

$$
\left\langle(\Delta n)^{2}\right\rangle=\left\langle n^{2}\right\rangle-\langle n\rangle^{2}=\langle n\rangle+\zeta^{2} T^{2}\left(\left\langle\bar{I}(t, T)^{2}\right\rangle-\bar{I}^{2}\right)
$$

Line width $\gamma\left(=1 / \tau_{c}\right) \quad$ Lorentzian lineshape function
Its second order coherence function is,

$$
g^{(2)}(\tau)=1+e^{-2 \gamma|\tau|}
$$

The average of integration time T is

$$
\begin{aligned}
& \left\langle g^{(2)}(0)\right\rangle-1=\frac{1}{T^{2}} \int_{0}^{T} \int_{0}^{T} e^{-2 \gamma \mid t_{2}-t_{1}} d t_{1} d t_{2} \\
& =\frac{1}{2 \gamma^{2} T^{2}}\left[e^{-2 \gamma T}+2 \gamma T-1\right]
\end{aligned}
$$

$$
\left\langle(\Delta n)^{2}\right\rangle=\langle n\rangle+\frac{\langle n\rangle^{2}}{2 \gamma^{2} T^{2}}\left[e^{-2 \gamma T}+2 \gamma T-1\right]
$$

$$
T \ll \tau_{c}
$$

$$
\left\langle(\Delta n)^{2}\right\rangle=\langle n\rangle+\langle n\rangle^{2}
$$

$$
\begin{aligned}
& T>\tau_{c} \\
& \left\langle(\Delta n)^{2}\right\rangle=\langle n\rangle+\frac{\langle n\rangle^{2}}{\gamma T}
\end{aligned}
$$

$T \gg \tau_{c}$
$\left\langle(\Delta n)^{2}\right\rangle=\langle n\rangle$

Measurement time : T
Redline: pulse



Variance of $\left\langle(\Delta n)^{2}\right\rangle$ by the velocity $v$


| Input <br> Voltage $(\mathrm{V})$ | frequency <br> $\left(\mathrm{s}^{-1}\right)$ | $\omega(=2 \pi \mathrm{f})$ <br> $\left(\mathrm{s}^{-1}\right)$ | $v(=r \omega)$ <br> $(\mathrm{m} / \mathrm{s})$ |
| :---: | :---: | :---: | :---: |
| 2 | 1.026 | 6.45 | 0.26 |
| 5 | 3.226 | 20.27 | 0.81 |
| 10 | 6.873 | 43.18 | 1.73 |
| 15 | 10.526 | 66.14 | 2.65 |
| 20 | 14.286 | 89.76 | 3.59 |
| 25 | 17.391 | 109.27 | 4.37 |

$$
\tau_{c}=\frac{\sigma(=9 \mu m)}{v} \quad\left\langle(\Delta n)^{2}\right\rangle=\langle n\rangle+\frac{\tau_{c}\langle n\rangle^{2}}{T}
$$

| Input Voltage | $\sigma / v$ 에 의해 <br> 계산된 $\tau_{r}$ | 광전자 분포에 의해 <br> 측정된 $\tau$ |
| :---: | :---: | :---: |
| $\mathrm{T}=50 \mu \mathrm{~s} 2 \mathrm{~V}$ | $34.62 \mu \mathrm{~s}$ | $31.4 \mu \mathrm{~s}$ |
| $\mathrm{~T}=30 \mu \mathrm{~s} 5 \mathrm{~V}$ | $11.11 \mu \mathrm{~s}$ | $11.4 \mu \mathrm{~s}$ |
| 10 V | $5.202 \mu \mathrm{~s}$ | $5.38 \mu \mathrm{~s}$ |
| $\mathrm{~T}=20 \mu \mathrm{~V}$ | $3.396 \mu \mathrm{~s}$ | $3.67 \mu \mathrm{~s}$ |
| 15 V | $2.507 \mu \mathrm{~s}$ | $2.92 \mu \mathrm{~s}$ |
| 20 V | $2.059 \mu \mathrm{~s}$ | $2.48 \mu \mathrm{~s}$ |

